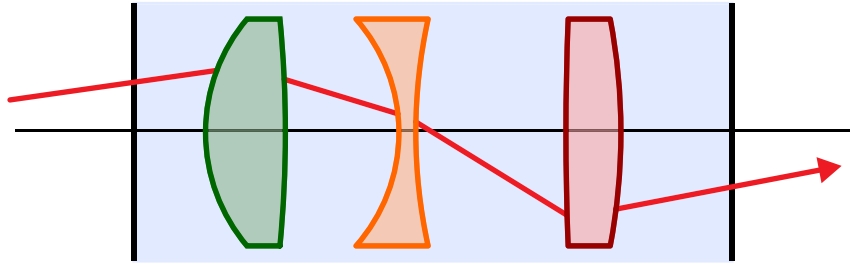


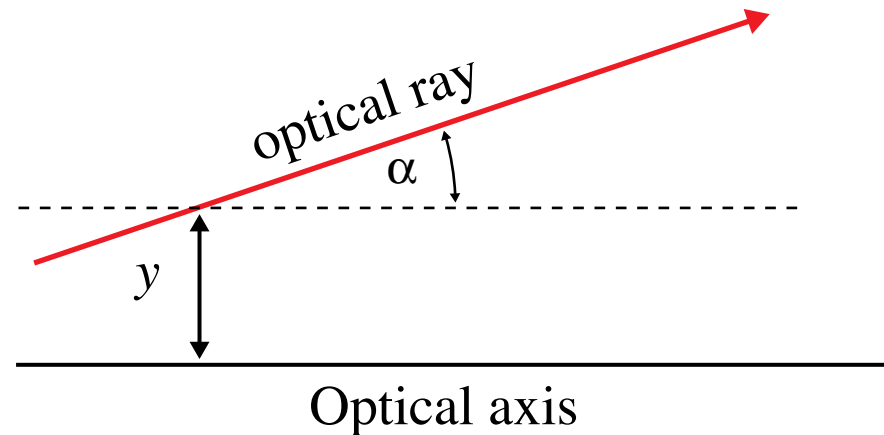
The Ray Vector



A light ray can be defined by two coordinates:

its “height”, y

its slope, α



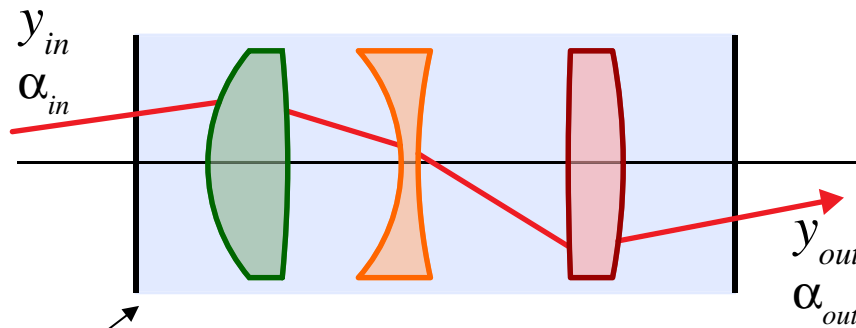
These parameters define a **ray vector**, which will change with distance, and/or as the ray propagates through optical interfaces and elements.

$$\begin{bmatrix} y \\ \alpha \end{bmatrix}$$

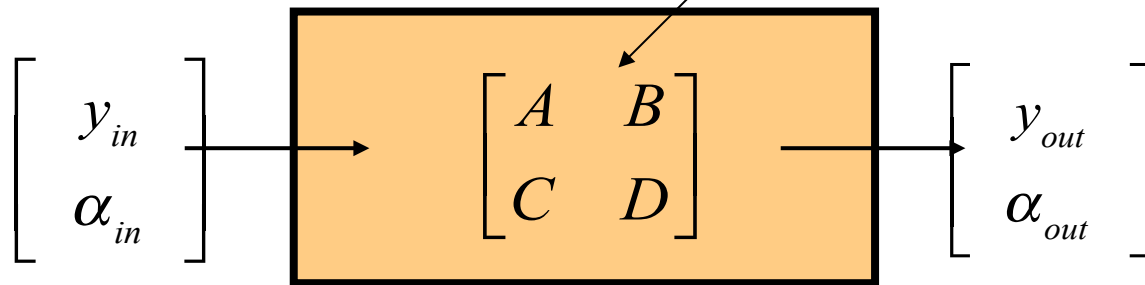
Ray Matrices

For many optical components, we can define 2 x 2 **ray matrices**.

An element's affect on a ray is found by multiplying its ray vector.



Optical system \leftrightarrow 2 x 2 Ray matrix



Ray matrices can describe simple and complex systems.

These matrices are often called ABCD Matrices.

Ray matrices as derivatives

Since the displacements and angles are assumed to be small, we can think in terms of partial derivatives.

$$y_{out} = \frac{\partial y_{out}}{\partial y_{in}} y_{in} + \frac{\partial y_{out}}{\partial \alpha_{in}} \alpha_{in}$$

$$\alpha_{out} = \frac{\partial \alpha_{out}}{\partial y_{in}} y_{in} + \frac{\partial \alpha_{out}}{\partial \alpha_{in}} \alpha_{in}$$

Diagram illustrating the matrix representation of the ray equations. The output vector $\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix}$ is equal to a matrix $\begin{bmatrix} A & B \\ C & D \end{bmatrix}$ multiplied by the input vector $\begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$. The matrix elements are labeled with partial derivatives:

- $A = \frac{\partial y_{out}}{\partial y_{in}}$ (spatial magnification)
- $B = \frac{\partial y_{out}}{\partial \alpha_{in}}$
- $C = \frac{\partial \alpha_{out}}{\partial y_{in}}$
- $D = \frac{\partial \alpha_{out}}{\partial \alpha_{in}}$ (angular magnification)

We can write these equations in matrix form.

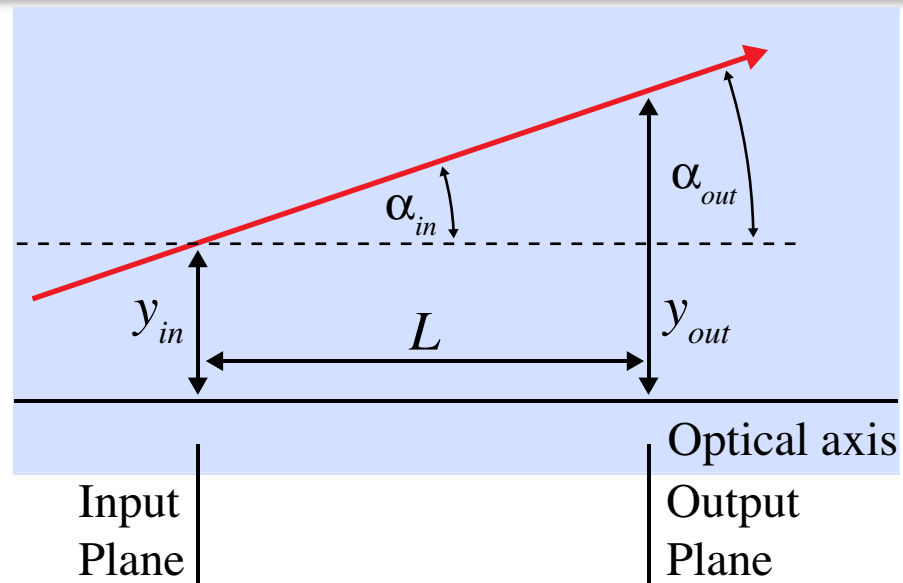
For cascaded elements, we simply multiply ray matrices.



$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = M_3 \left\{ M_2 \left(M_1 \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix} \right) \right\} = M_3 M_2 M_1 \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

Notice that the order looks opposite to what you think it should be, but it makes sense when you think about it.

Translation Matrix



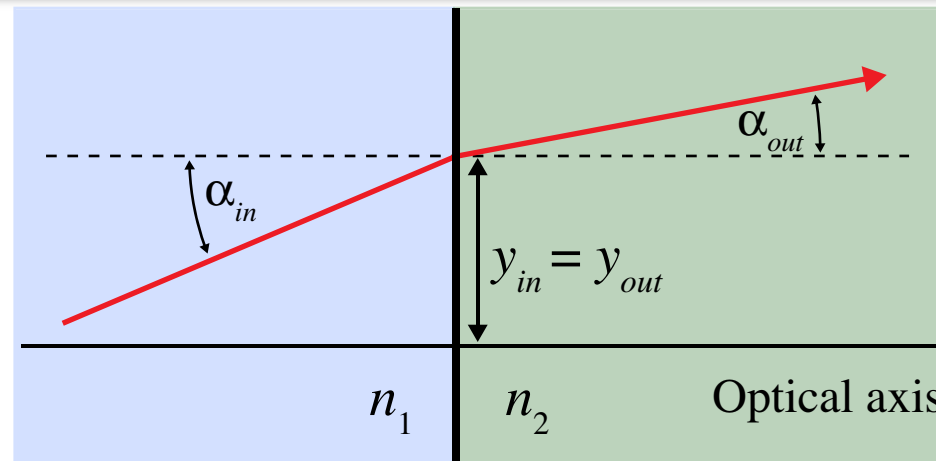
$$\alpha_{out} = \alpha_{in} \quad y_{out} = y_{in} + L \tan \alpha_{in} \cong y_{in} + L \alpha_{in}$$

$$y_{out} = (1)y_{in} + (L)\alpha_{in}$$

$$\alpha_{out} = (0)y_{in} + (1)\alpha_{in}$$

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix} = \begin{bmatrix} 1 & x_{out} - x_{in} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

Refraction Matrix – Flat Interface



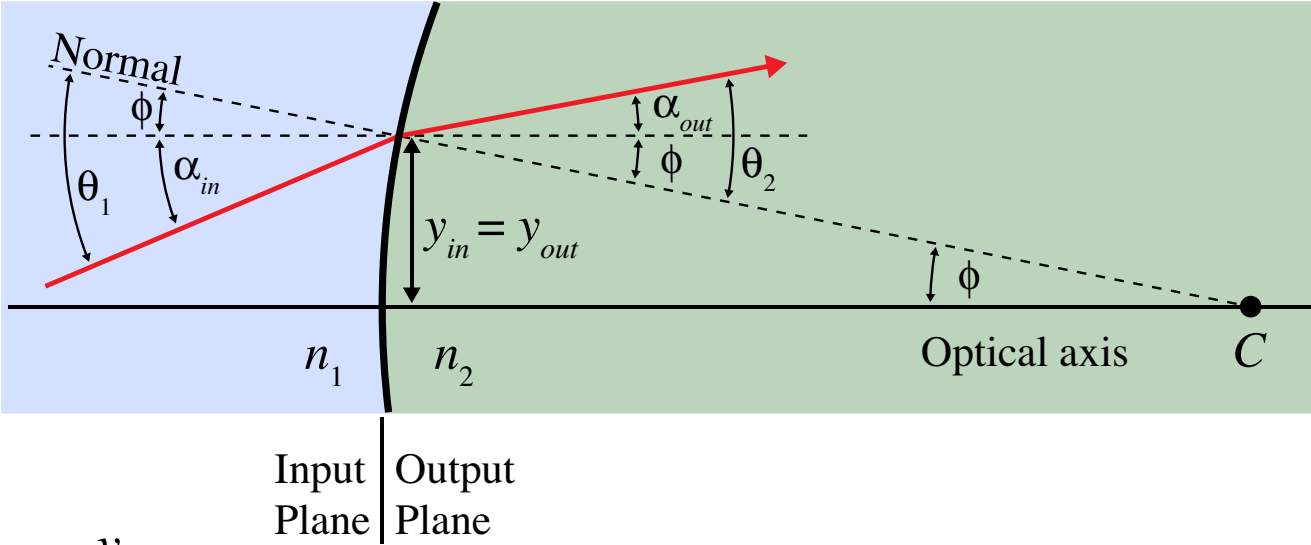
$$y_{out} = (1)y_{in} + (0)\alpha_{in}$$

Input Plane	Output Plane
-------------	--------------

Paraxial Snell's Law: $n_1 \alpha_{in} = n_2 \alpha_{out} \Rightarrow \alpha_{out} = (0)y_{in} + \left(\frac{n_1}{n_2}\right)\alpha_{in}$

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & \left(\frac{n_1}{n_2}\right) \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

Refraction Matrix – Curved Interface



$$\alpha_{out} = \theta_2 - \phi = \theta_2 - \frac{y_{in}}{R}$$

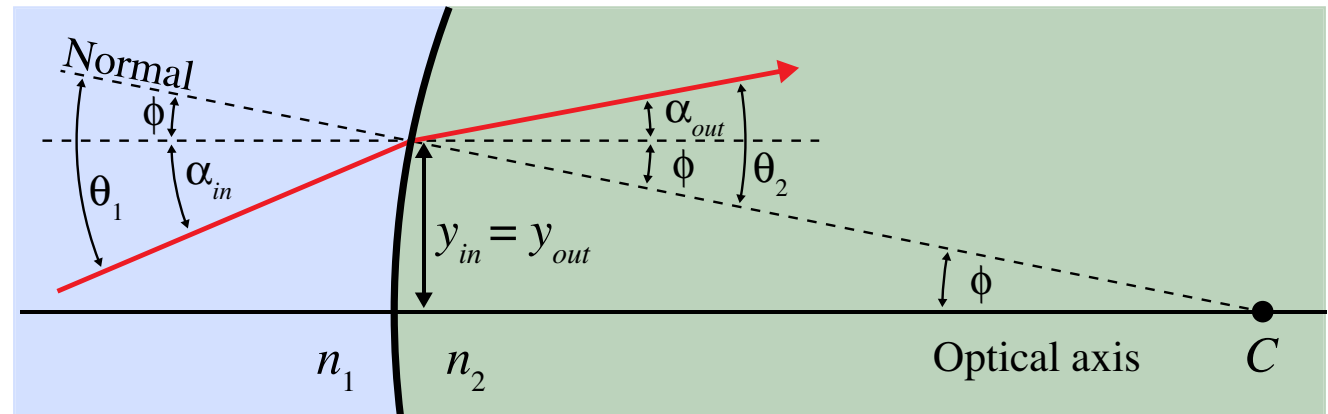
$$\alpha_{in} = \theta_1 - \phi = \theta_1 - \frac{y_{in}}{R}$$

$$\theta_1 = \alpha_{in} + \frac{y_{in}}{R} \quad \theta_2 = \alpha_{out} + \frac{y_{in}}{R}$$

Paraxial Snell's Law: $n_1 \theta_1 = n_2 \theta_2 \rightarrow n_1 \left(\alpha_{in} + \frac{y_{in}}{R} \right) = n_2 \left(\alpha_{out} + \frac{y_{in}}{R} \right)$

$$\alpha_{out} = \left(\frac{n_1}{n_2} \right) \left(\alpha_{in} + \frac{y_{in}}{R} \right) - \frac{y_{in}}{R} = \boxed{ \frac{1}{R} \left(\frac{n_1}{n_2} - 1 \right) y_{in} + \left(\frac{n_1}{n_2} \right) \alpha_{in} }$$

Refraction Matrix – Curved Interface



$$y_{out} = (1)y_{in} + (0)\alpha_{in}$$

Input Plane | Output Plane

$$\alpha_{out} = \left[\frac{1}{R} \left(\frac{n_1}{n_2} - 1 \right) \right] y_{in} + \left(\frac{n_1}{n_2} \right) \alpha_{in}$$

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ \frac{1}{R} \left(\frac{n_1}{n_2} - 1 \right) & \left(\frac{n_1}{n_2} \right) \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

Concave surface: $R < 0$

Convex surface: $R > 0$

Reflection Matrix

$$\alpha_{out} = \theta_1 + \theta_2 - \alpha_{in} \quad \theta_1 = \alpha_{in} - \phi = \alpha_{in} - \frac{y_{in}}{-R} = \alpha_{in} + \frac{y_{in}}{R}$$

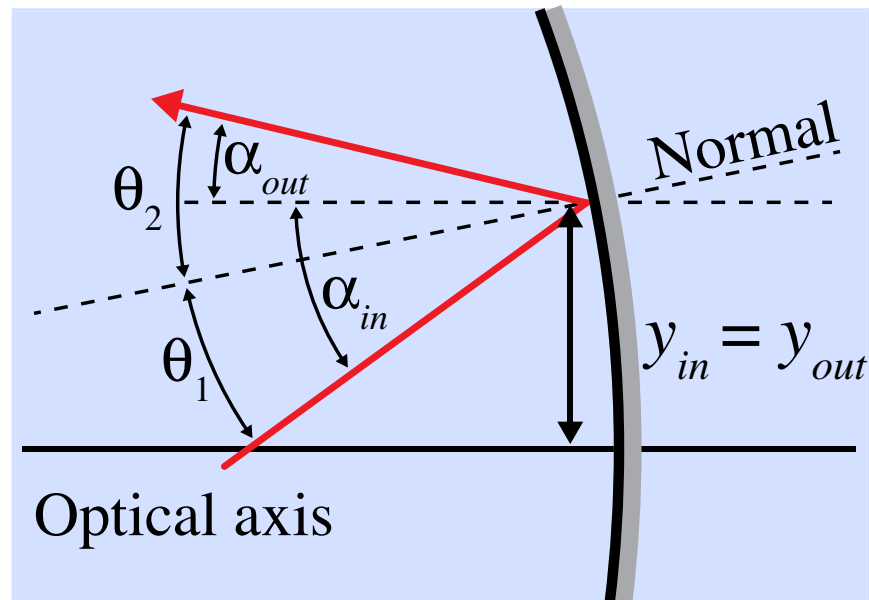
Law of Reflection: $\theta_1 = \theta_2$

$$\alpha_{out} = 2\theta_1 - \alpha_{in} = \alpha_{in} + \frac{2}{R}y_{in}$$

$$y_{out} = (1)y_{in} + (0)\alpha_{in}$$

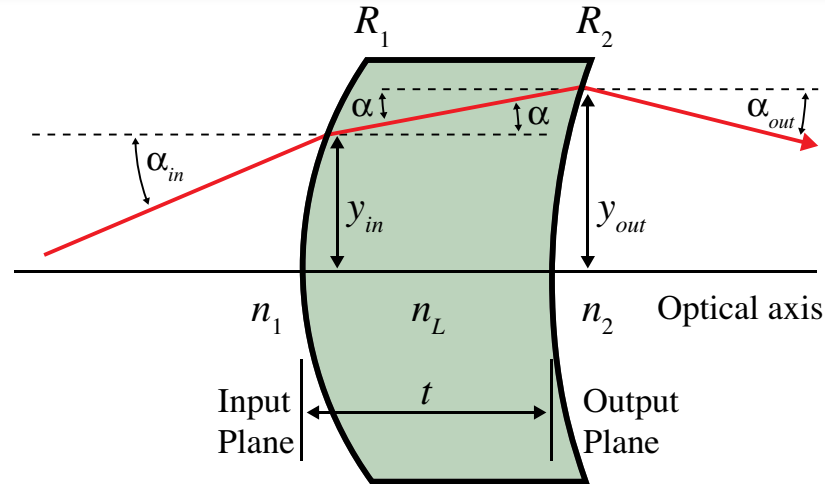
$$\alpha_{out} = \left(\frac{2}{R}\right)y_{in} + (1)\alpha_{in}$$

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ \frac{2}{R} & 1 \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$



Input Plane | Output Plane

Thick Lens Matrix



Refraction at first surface:

$$\begin{bmatrix} y_1 \\ \alpha \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ \frac{1}{R_1} \left(\frac{n_1}{n_L} - 1 \right) & \frac{n_1}{n_L} \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix} = M_1 \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

Translation from 1st surface to 2nd surface:

$$\begin{bmatrix} y_2 \\ \alpha \end{bmatrix} = \begin{bmatrix} 1 & t \\ 0 & 1 \end{bmatrix} \begin{bmatrix} y_1 \\ \alpha \end{bmatrix} = M_2 \begin{bmatrix} y_1 \\ \alpha \end{bmatrix}$$

Refraction at second surface:

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ \frac{1}{R_2} \left(\frac{n_L}{n_2} - 1 \right) & \frac{n_L}{n_2} \end{bmatrix} \begin{bmatrix} y_2 \\ \alpha \end{bmatrix} = M_3 \begin{bmatrix} y_2 \\ \alpha \end{bmatrix}$$

Thick Lens Matrix

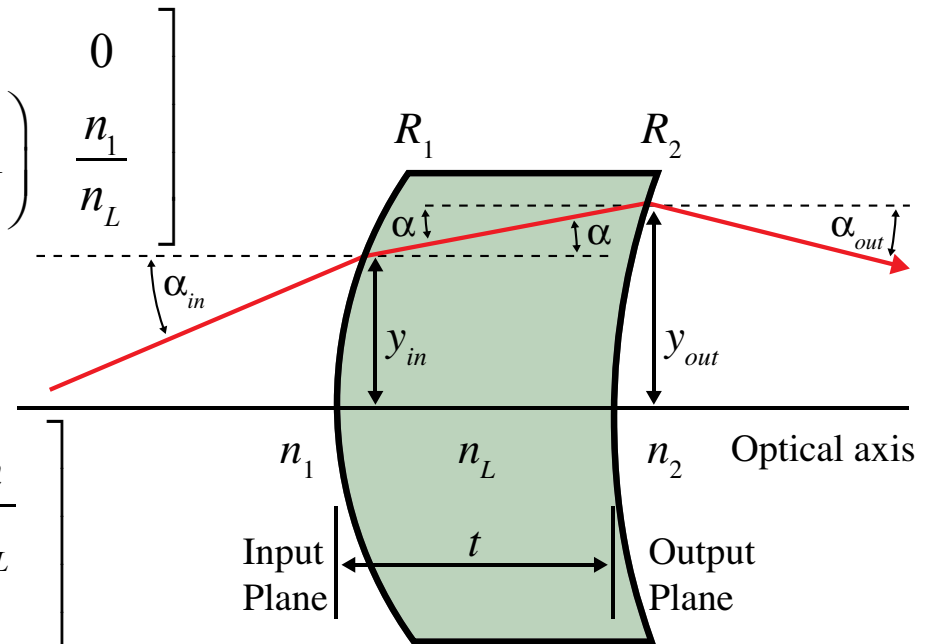
Thick lens matrix: $M = M_3 M_2 M_1$

$$M = \begin{bmatrix} 1 & 0 \\ \frac{1}{R_2} \left(\frac{n_L}{n_2} - 1 \right) & \frac{n_L}{n_2} \end{bmatrix} \begin{bmatrix} 1 & t \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ \frac{1}{R_1} \left(\frac{n_1}{n_L} - 1 \right) & \frac{n_1}{n_L} \end{bmatrix}$$

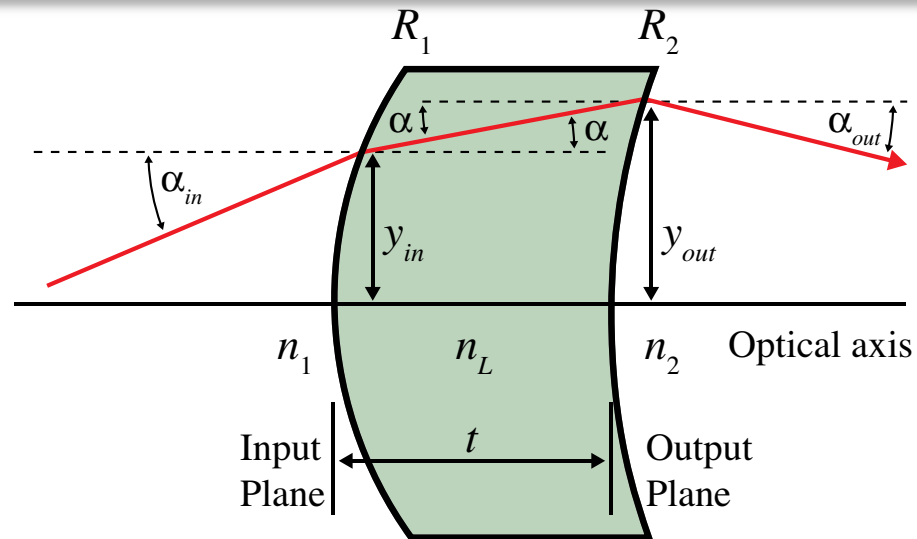
Assuming $n_1 = n_2 = n$:

$$M = \begin{bmatrix} 1 & 0 \\ \frac{1}{R_2} \left(\frac{n_L}{n} - 1 \right) & \frac{n_L}{n} \end{bmatrix} \begin{bmatrix} 1 + \frac{t}{R_1} \left(\frac{n}{n_L} - 1 \right) & t \frac{n}{n_L} \\ \frac{1}{R_1} \left(\frac{n}{n_L} - 1 \right) & \frac{n}{n_L} \end{bmatrix}$$

$$= \begin{bmatrix} 1 + \frac{t}{R_1} \left(\frac{n}{n_L} - 1 \right) & t \frac{n}{n_L} \\ \frac{1}{R_2} \left(\frac{n_L}{n} - 1 \right) \left[1 + \frac{t}{R_1} \left(\frac{n}{n_L} - 1 \right) \right] - \frac{1}{R_1} \left(\frac{n_L}{n} - 1 \right) & 1 - \frac{t}{R_2} \left(\frac{n}{n_L} - 1 \right) \end{bmatrix}$$



Thick Lens Matrix



$$M = \begin{bmatrix} 1 + \frac{t}{R_1} \left(\frac{n}{n_L} - 1 \right) & t \frac{n}{n_L} \\ \frac{1}{R_2} \left(\frac{n_L}{n} - 1 \right) \left[1 + \frac{t}{R_1} \left(\frac{n}{n_L} - 1 \right) \right] - \frac{1}{R_1} \left(\frac{n_L}{n} - 1 \right) & 1 - \frac{t}{R_2} \left(\frac{n}{n_L} - 1 \right) \end{bmatrix}$$

Thin Lens Matrix

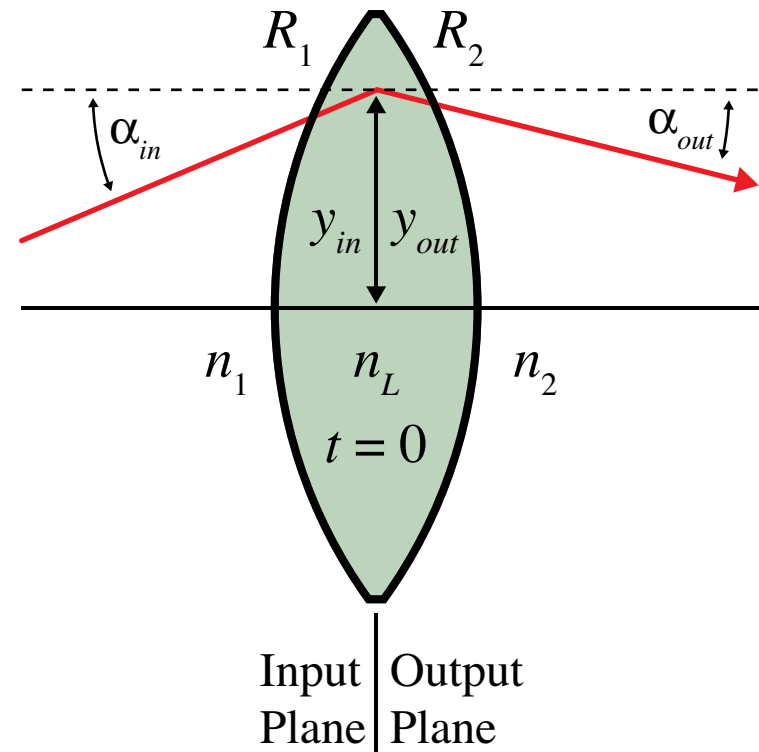
The thin lens matrix is found by setting $t = 0$:

Thin lens matrix:

$$M = \begin{bmatrix} 1 & 0 \\ \left(\frac{n_L}{n} - 1\right) \left(\frac{1}{R_2} - \frac{1}{R_1}\right) & 1 \end{bmatrix}$$

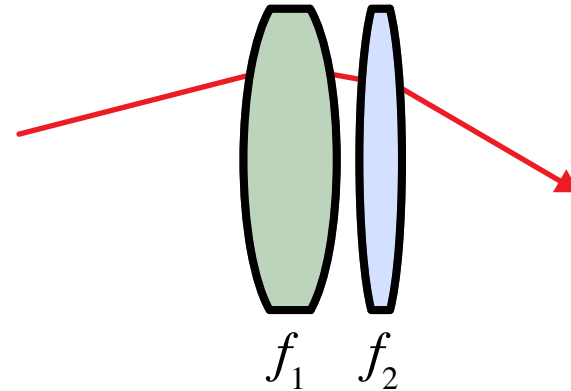
but $\frac{1}{f} = \left(\frac{n_L}{n} - 1\right) \left(\frac{1}{R_1} - \frac{1}{R_2}\right)$

$$M = \begin{bmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{bmatrix}$$



Consecutive thin lenses

Suppose we have two lenses right next to each other (with no space in between).



$$M_{tot} = \begin{bmatrix} 1 & 0 \\ -1/f_2 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -1/f_1 & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ -1/f_1 - 1/f_2 & 1 \end{bmatrix}$$
$$1/f_{tot} = 1/f_1 + 1/f_2$$

So two consecutive lenses act as one whose focal length is computed by the resistive sum.

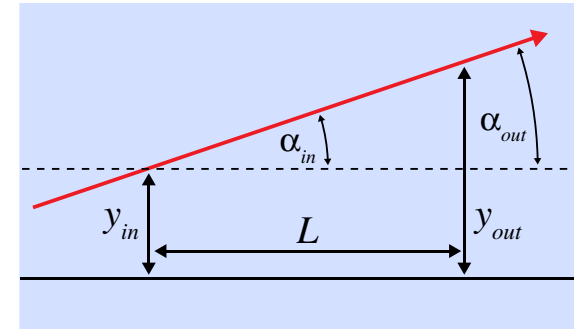
As a result, we define a measure of **inverse** lens focal length, the **dioptr**.

$$1 \text{ dioptr} = 1 \text{ m}^{-1}$$

Summary of Matrix Methods

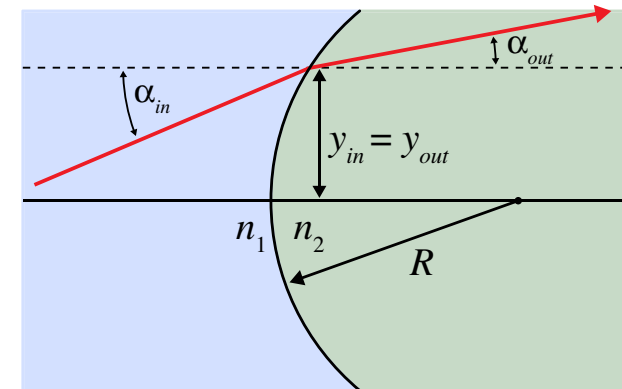
Translation Matrix:

$$M = \begin{bmatrix} 1 & L \\ 0 & 1 \end{bmatrix}$$



Refraction matrix,
spherical interface:

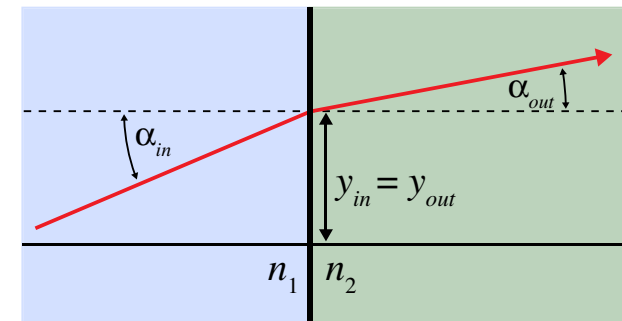
$$M = \begin{bmatrix} 1 & 0 \\ \frac{1}{R} \left(\frac{n_1}{n_2} - 1 \right) & \frac{n_1}{n_2} \end{bmatrix}$$



(+R): Convex (-R): Concave

Refraction matrix,
plane interface:

$$M = \begin{bmatrix} 1 & 0 \\ 0 & \frac{n_1}{n_2} \end{bmatrix}$$

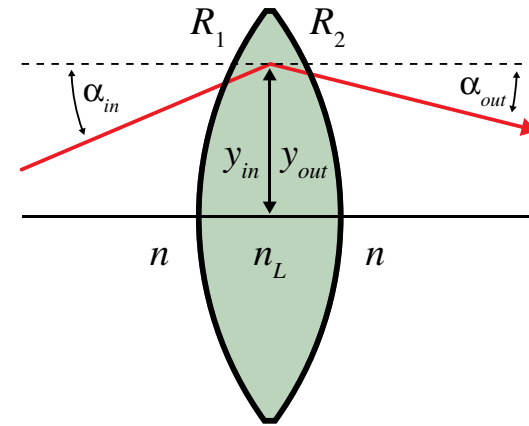


Summary of Matrix Methods

Thin-lens Matrix:

$$M = \begin{bmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{bmatrix}$$

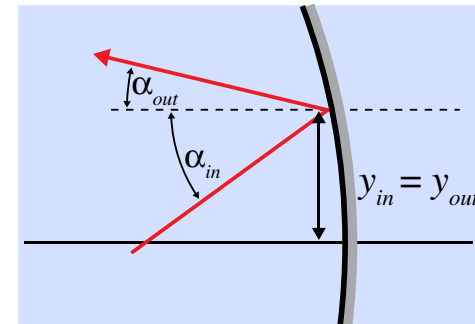
$$\frac{1}{f} = \left(\frac{n_L}{n} - 1\right) \left(\frac{1}{R_1} - \frac{1}{R_2}\right)$$



(+f): Converging (-f): Diverging

Spherical mirror
matrix:

$$M = \begin{bmatrix} 1 & 0 \\ \frac{2}{R} & 1 \end{bmatrix}$$



(+R): Convex (-R): Concave

System Ray-Transfer Matrix

Any paraxial optical system, no matter how complicated, can be represented by a 2x2 optical matrix. This matrix M is usually denoted

$$M = \begin{bmatrix} A & B \\ C & D \end{bmatrix}$$

A useful property of this matrix is that

$$\text{Det } M = AD - BC = \frac{n_{in}}{n_{out}}$$

where n_{in} and n_{out} are the refractive indices of the input and output media of the optical system. Usually, the medium will be air on both sides of the optical system and

$$\text{Det } M = AD - BC = \frac{n_{in}}{n_{out}} = 1$$

System Ray-Transfer Matrix

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix} \Rightarrow$$

$$y_{out} = Ay_{in} + B\alpha_{in}$$

$$\alpha_{out} = Cy_{in} + D\alpha_{in}$$

Let's examine the implications when any of the four elements of the system matrix is equal to zero.

System Ray-Transfer Matrix

Let's see what happens when $D = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & B \\ C & 0 \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = Ay_{in} + B\alpha_{in}$$

$$\alpha_{out} = Cy_{in}$$

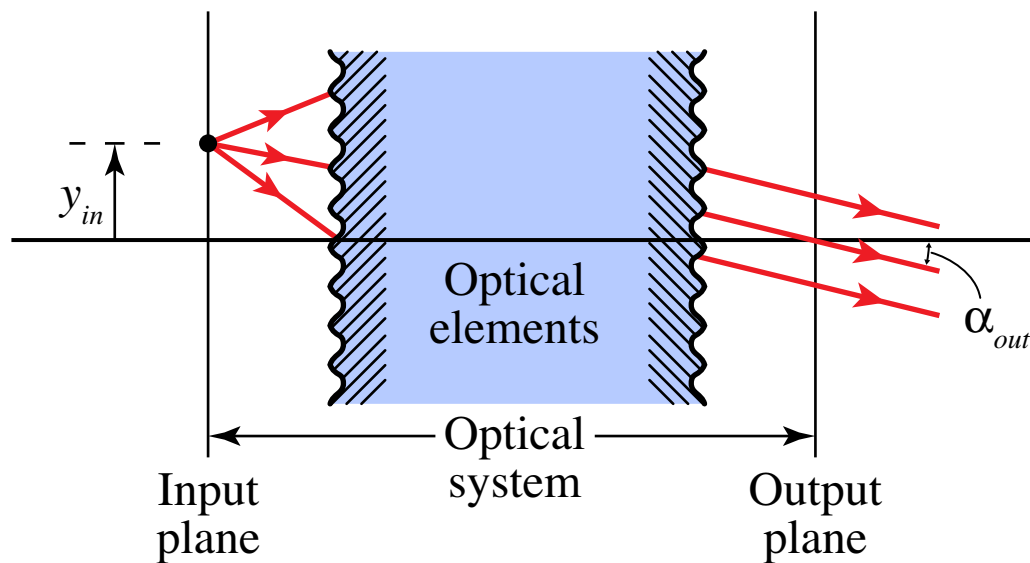
System Ray-Transfer Matrix

Let's see what happens when $D = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & B \\ C & 0 \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = Ay_{in} + B\alpha_{in}$$

$$\alpha_{out} = Cy_{in}$$



When $D = 0$, the input plane for the optical system is the input focal plane.

System Ray-Transfer Matrix

Let's see what happens when $A = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 0 & B \\ C & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = B\alpha_{in}$$

$$\alpha_{out} = Cy_{in} + D\alpha_{in}$$

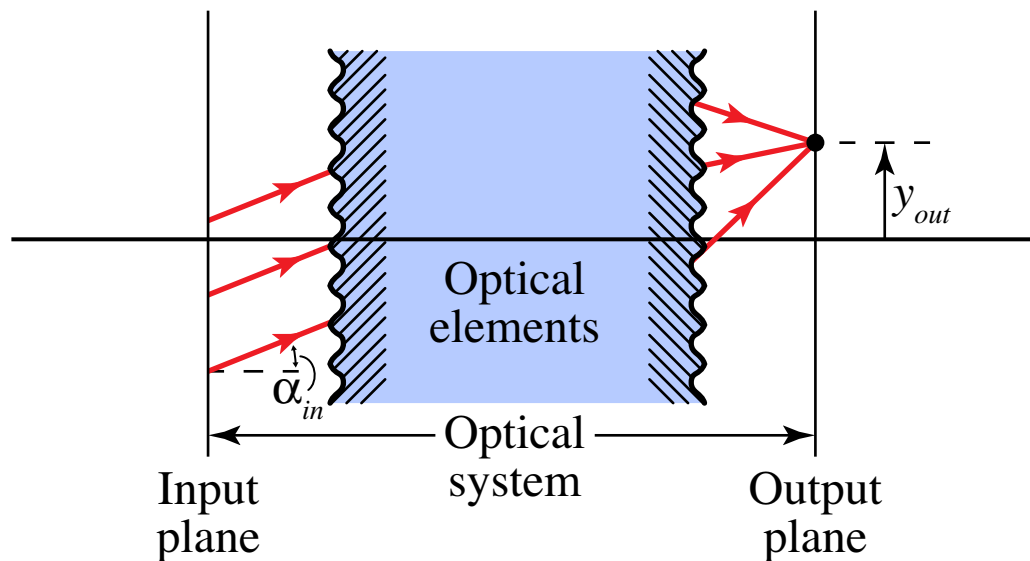
System Ray-Transfer Matrix

Let's see what happens when $A = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} 0 & B \\ C & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = B\alpha_{in}$$

$$\alpha_{out} = Cy_{in} + D\alpha_{in}$$



When $A = 0$, the output plane for the optical system is the output focal plane.

System Ray-Transfer Matrix

Let's see what happens when $C = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & B \\ 0 & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = Ay_{in} + B\alpha_{in}$$

$$\alpha_{out} = D\alpha_{in}$$

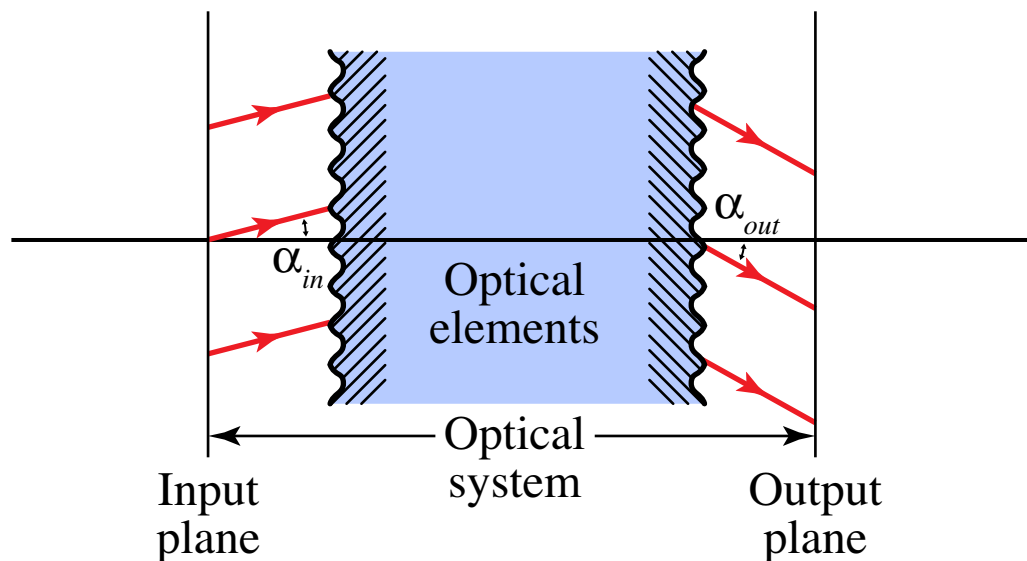
System Ray-Transfer Matrix

Let's see what happens when $C = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & B \\ 0 & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = Ay_{in} + B\alpha_{in}$$

$$\alpha_{out} = D\alpha_{in}$$



When $C = 0$, collimated light at the input plane is collimated light at the output plane, but the angle with the optical axis is different. This is a telescopic arrangement, with an angular magnification of $D = \alpha_{out} / \alpha_{in}$.

System Ray-Transfer Matrix

Let's see what happens when $B = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & 0 \\ C & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = Ay_{in}$$

$$\alpha_{out} = Cy_{in} + D\alpha_{in}$$

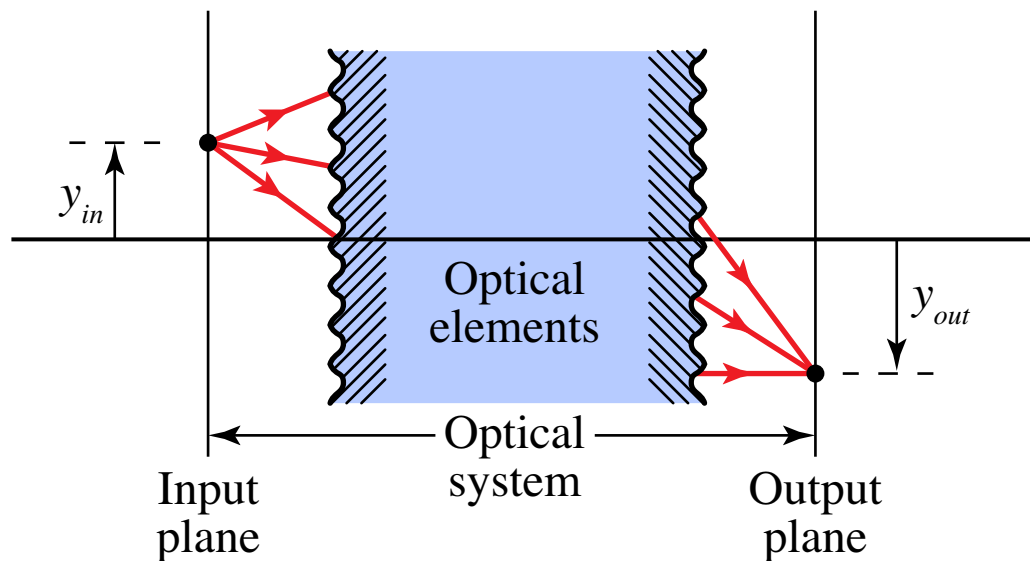
System Ray-Transfer Matrix

Let's see what happens when $B = 0$.

$$\begin{bmatrix} y_{out} \\ \alpha_{out} \end{bmatrix} = \begin{bmatrix} A & 0 \\ C & D \end{bmatrix} \begin{bmatrix} y_{in} \\ \alpha_{in} \end{bmatrix}$$

$$y_{out} = Ay_{in}$$

$$\alpha_{out} = Cy_{in} + D\alpha_{in}$$



When $B = 0$, the input and output planes are object and image planes, respectively, and the transverse magnification of the system $m = A$.

System Ray-Transfer Matrix – Summary

When $D = 0$, the input plane for the optical system is the input focal plane.

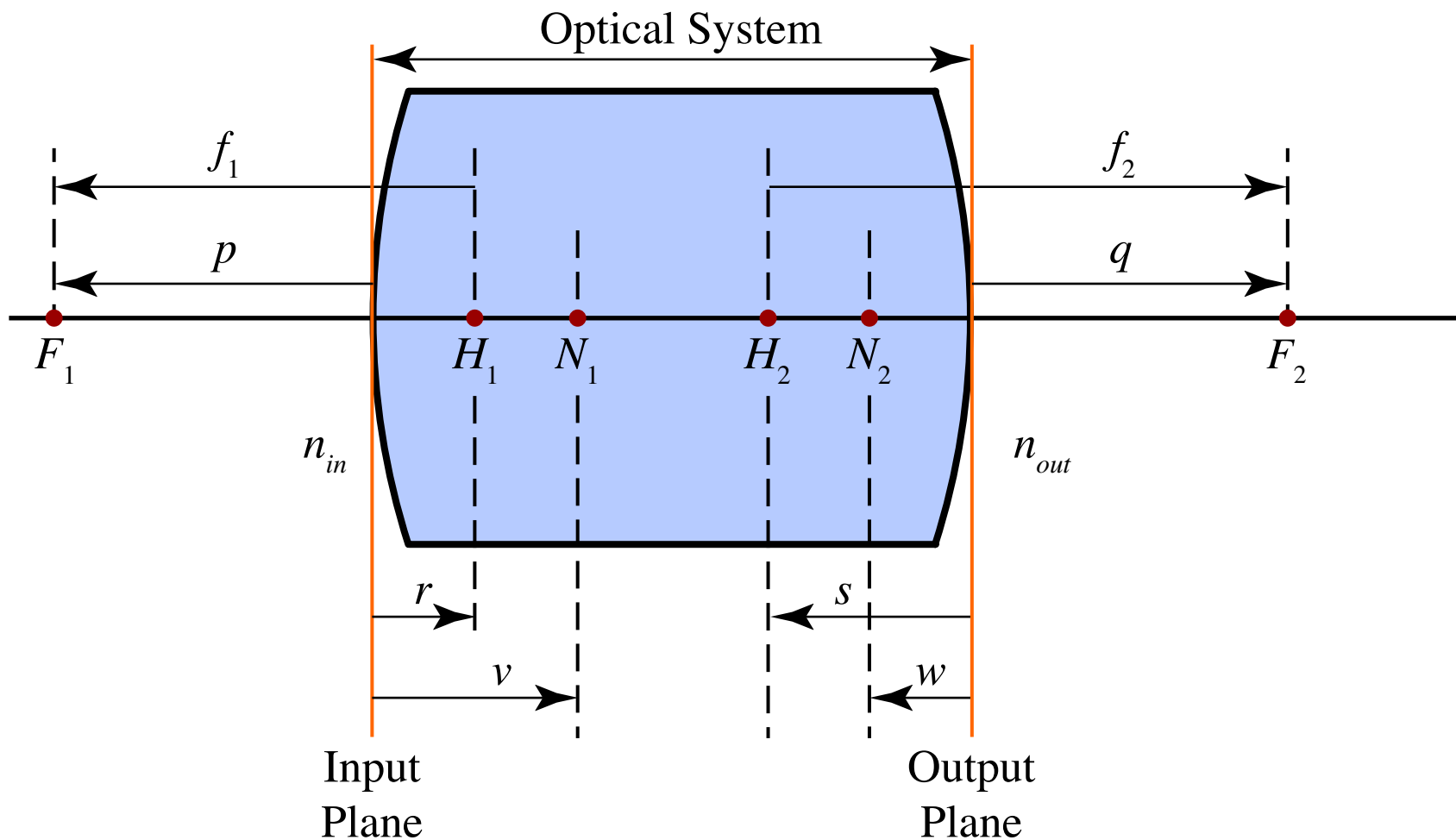
When $A = 0$, the output plane for the optical system is the output focal plane.

When $B = 0$, the input and output planes are object and image planes, respectively, and the transverse magnification of the system $m = A$.

When $C = 0$, collimated light at the input plane is collimated light at the output plane, but the angle with the optical axis is different. This is a telescopic arrangement, with an angular magnification of $D = \alpha_{out} / \alpha_{in}$.

System Ray-Transfer Matrix

The matrix elements of the system matrix can be analyzed to determine the cardinal points and planes of an optical system.



Cardinal Points

$$p = \frac{D}{C}$$

 F_1

$$q = -\frac{A}{C}$$

 F_2

$$r = \frac{D - n_{in}/n_{out}}{C}$$

 H_1

$$s = \frac{1 - A}{C}$$

 H_2

$$v = \frac{D - 1}{C}$$

 N_1

$$w = \frac{n_{in}/n_{out} - A}{C}$$

 N_2

Located relative to input (1) and output (2) reference planes.

+ values are to the right of the plane.
- values are to the left of the plane.

$$f_1 = p - r = \frac{n_{in}/n_{out}}{C}$$

 F_1

$$f_2 = q - s = -\frac{1}{C}$$

 F_2

Located relative to principal planes.